

# Blackbody radiation

- Note: I am not following the derivation directly from the website  
<http://www.cv.nrao.edu/course/astr534/PDFnew.shtml>
- It needs QM treatment to understand where it comes from (classical treatment has no limit to power you would emit). I use one from Harwit
  - *Martin Harwit 'Astrophysical Concepts'*
- You **do not need to repeat** the derivation, but do need to understand it!

# 6D Phase space

- In quantum mechanics particles are identical if they have the same spin and occupy the same cell in phase space (3D position, 3D momentum)

$$\delta x \delta y \delta z \delta p_x \delta p_y \delta p_z = h^3$$

- so if we allow up to a maximum momentum  $p_{\max}$  and 2 spin states we have a max number of particles in a volume  $V$ :

$$\frac{8 \pi p_{\max}^3 V}{3 h^3}$$

# Photons

- 2 spins (LHC, RHC) bosons

- Freq  $\nu$

$$\nu \equiv \frac{pc}{h}$$

- number of phase space cells from  $p$  to  $p+dp$

$$Z(p) dp = 2V \frac{4\pi p^2 dp}{h^3}$$

- so from  $\nu$  to  $\nu+d\nu$

$$Z(\nu) = 2 \left[ \frac{4\pi \nu^2 d\nu}{c^3} \right] V$$

# States & Probabilities

- Bosons can aggregate with zero point energy
- Relative probability for photon to be in the  $n^{\text{th}}$  energy state at temperature  $T$   
 $e^{-((n+\frac{1}{2})h\nu)/kt}$
- Absolute probability

$$\frac{e^{-((n+\frac{1}{2})h\nu)/kt}}{\sum_n e^{-((n+\frac{1}{2})h\nu)/kt}} = \frac{e^{-nh\nu/kt}}{\sum_n e^{-nh\nu/kt}}$$

# Average energy per phase cell

- energy of that cell times its probability

$$\frac{\sum_n (n + \frac{1}{2}) h\nu e^{-nh\nu/kt}}{\sum_n e^{-nh\nu/kt}}$$

- replace  $h\nu/kT$  by  $x$

$$\langle E \rangle = \frac{kT (xe^{-x} + 2xe^{-2x} + 3xe^{-3x} + \dots)}{1 + e^{-x} + e^{-2x} + e^{-3x} + \dots} + h\nu/2$$

# continued..

- using  $(1+a+a^2+a^3\dots) = 1/(1-a)$
- numerator

$$kT [(x(e^{-x} + e^{-2x} + e^{-3x} + \dots)) + x(e^{-2x} + e^{-3x} + \dots) + x(e^{-3x\dots}) + \dots]$$

$$kT \left[ \frac{xe^{-x}}{1-e^{-x}} + \frac{xe^{-2x}}{1-e^{-x}} + \frac{xe^{-3x}}{1-e^{-x}} + \dots \right]$$

$$\frac{kTxe^{-x}}{(1-e^{-x})^2}$$

- denominator

$$\frac{1}{(1-e^{-x})}$$

# Finally

- combining back

$$\langle E \rangle = \frac{kTx e^{-x}}{1 - e^{-x}} + \frac{h\nu}{2} = \frac{kTx}{e^x - 1} + \frac{h\nu}{2} = \frac{h\nu}{e^{h\nu/kT} - 1} + \frac{h\nu}{2}$$

- so energy density per phase cell times phase cells per unit volume gives

$$\rho(\nu, T) d\nu = \frac{8\pi\nu^2 d\nu}{c^3} \left( \frac{h\nu}{e^{h\nu/kT} - 1} + \frac{h\nu}{2} \right)$$

- we can ignore  $\frac{1}{2}h\nu$  as it is not observable in emission or absorption; it is a zero-point

# Blackbody

- General form

$$\rho(\nu, T) d\nu = \frac{8\pi\nu^2 d\nu}{c^3} \left( \frac{h\nu}{e^{h\nu/kT} - 1} \right)$$

- Intensity per unit area per unit frequency per unit solid angle (expansion at  $c$  over  $4\pi$  steradians)

$$I(\nu, T) = \frac{c\rho(\nu, T)}{4\pi} = \frac{2\nu^2}{c^2} \left( \frac{h\nu}{e^{h\nu/kT} - 1} \right)$$



# Stefan's Law

- integrate to infinity

$$\int_0^{\infty} \frac{2h\nu^3}{c^2} \left( \frac{1}{e^{h\nu/kT} - 1} \right) d\nu = \sigma T^4$$

- where  $\sigma$  is Stefan's constant, the radiation per unit area

$$\sigma = \frac{2\pi^5 k^4}{15h^3 c^2}$$

$$5.67 \times 10^{-8} \text{ Wm}^{-2} \text{ K}^{-4}$$

# Approximations

- At low frequency (radio)  $e^{h\nu/kT} \sim 1+h\nu/kT$

$$B(\nu, T) = \frac{c \rho(\nu, T)}{4\pi} = \frac{2\nu^2}{c^2} \left( \frac{h\nu}{e^{h\nu/kT} - 1} \right) \simeq \frac{2kT\nu^2}{c^2}$$

- At high frequency (optical and above)

$$e^{h\nu/kT} \gg 1$$

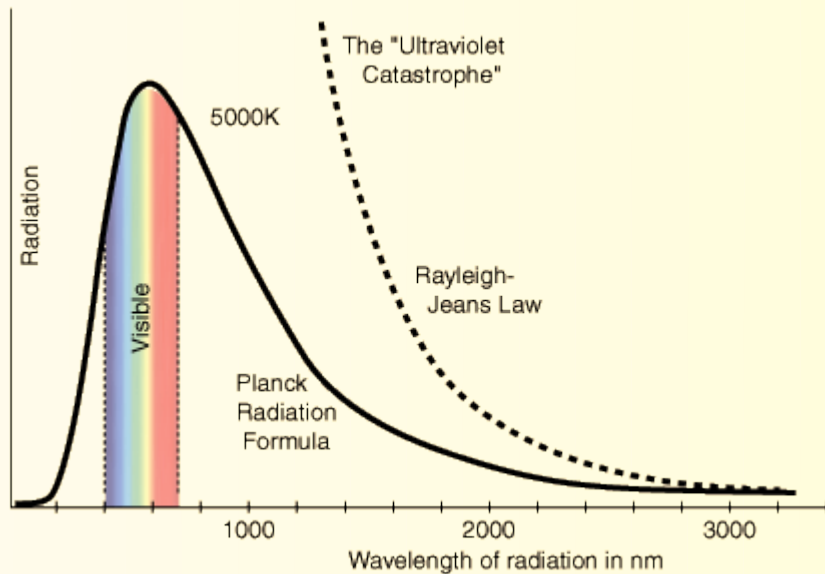
$$B(\nu, T) \simeq \frac{2h\nu^3}{c^2} \left( \frac{1}{e^{h\nu/kT}} \right)$$

# Wien's Law

- By differentiation

$$\frac{\partial B}{\partial \nu} = 0 \quad \text{when} \quad \nu = 59 \text{ GHz} \times T (K)$$

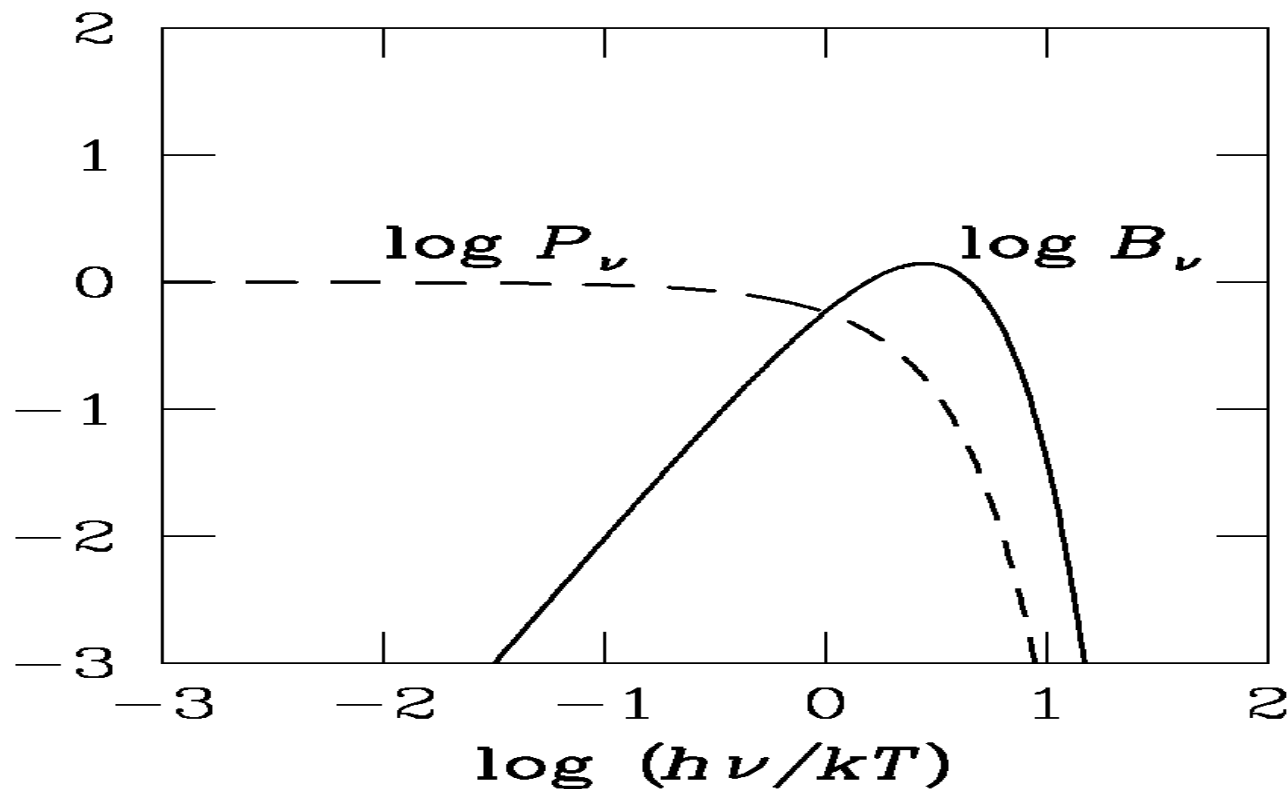
$$\frac{\partial B}{\partial \lambda} = 0 \quad \text{when} \quad \lambda = \frac{2.9 \text{ mm}}{T (K)}$$



# Power from warm resistor

- The resistor is the electrical equivalent of a black body in 1 dimension instead of 3

$$P(\nu, T) = \frac{h\nu}{e^{h\nu/kT} - 1} \simeq kT \text{ at low frequencies}$$



# Johnson-Nyquist Resistor noise

- Independent of resistance!
- Related to CCD dark current noise
- Depends on Temperature only at low freq.
- Comes from the Fluctuation-Dissipation theorem of statistical mechanics