

## Comptonization Spectra

The growth or evolution of a Compton spectrum is determined by the Kompaneets eqn.

$$\frac{\partial n}{\partial y} = \frac{1}{x^2} \frac{\partial}{\partial x} \left[ x^4 \left( \frac{\partial n}{\partial x} + n^2 + n \right) \right] \quad x = \frac{h\nu}{kT}$$

The LHS is

$$\frac{\partial n}{\partial y} = \frac{m_e c^2}{kT_e} \frac{1}{n_e \sigma_T c} \frac{\partial n}{\partial t}$$

$$\Rightarrow \frac{\partial}{\partial y} = \frac{m_e c^2}{kT} \frac{1}{n_e \sigma_T c} \frac{\partial}{\partial t}$$

$$\Rightarrow dy = \frac{kT_e}{m_e c^2} n_e \sigma_T c dt$$

$$= \frac{kT_e}{m_e c^2} n_e \sigma_T c dt$$

$$= \frac{kT_e}{m_e c^2} n_e \sigma_T dR$$

## The Sunyaev-Zeldovich effect

The S-Z effect is the distortion of the CMB background as a result of the IC scattering of hot electrons in the halos of galaxy clusters. We can calculate this effect by using the Kompaneets eqn. Because  $T_{\text{gas}}(\text{cluster}) \gg T_{\text{rad}}(\text{CMB})$  the  $\frac{dn}{dx}$  term in the Kompaneets eqn dominates. This is because

$$\left\{ \begin{array}{l} \text{Dimensional} \\ \text{analysis} \end{array} \right\} \quad \frac{dn}{dx} \approx \frac{n}{x} \gg n \quad \text{for} \quad x = \frac{h\nu}{kT} \ll 1$$

Then

$$\frac{dn}{dy} \approx \frac{1}{x^2} \frac{d}{dx} \left( x^4 \frac{dn}{dx} \right)$$

To quantify the perturbation on the CMB photon density as a result of IC scattering, we write

$$\delta n = \left[ \frac{dn}{dy} \right] dy$$

For a cluster

$$y = \int dy \\ = \int \frac{kT_e}{m_e c^2} n_e \sigma_T dR \ll 1 \\ \uparrow \\ \text{low}$$

$$y \rightarrow dy$$

$$\delta n = \frac{dn}{dy} y = y \left[ \frac{1}{x^2} \frac{d}{dx} \left( x^4 \frac{dn}{dx} \right) \right]$$

Therefore

$$\frac{\delta h}{h} = \frac{y}{n} \frac{dn}{dy}$$
$$= \frac{y}{n} \left[ \frac{1}{x} \frac{d}{dx} \left( x^4 \frac{dn}{dx} \right) \right]$$

For a Planck spectrum (CMB)

$$n = \frac{1}{e^x - 1} \quad x = \frac{h\nu}{kT}$$

$$\frac{\delta h}{h} = y \left[ \frac{x e^x}{e^x - 1} \right] \left[ x \coth\left(\frac{x}{2}\right) - 4 \right] \quad \text{[Verify it]}$$

In the R-J limit  $x = \frac{h\nu}{kT} \ll 1$

Show that:

$$\frac{\delta h}{h} = -2y \quad \left[ y = \int \frac{kT_e}{m_e c^2} n_e g_T dV \right]$$

Hence CMB photons upon interaction with cluster electrons will be pumped up in energy through IC scattering. The resultant effect

$$\left( \frac{\delta h}{h} \right)_{\text{clus}} = -2y$$

$\Rightarrow$  [This depletion of CMB photons is then added to CMB at higher frequencies.]

For a BB spectrum

$$n_{ph} = 0.24 \left( \frac{k}{hc} \right)^3 T^3$$

$$\delta n_{ph} = 0.24 \left( \frac{k}{hc} \right)^3 3 T^2 \delta T$$

Therefore

$$\frac{\delta n_{ph}}{n_{ph}} = \frac{3}{T_{bb}} \delta T_{bb} \propto \frac{\delta T_{bb}}{T_{bb}}$$

Hence

$$\frac{\delta T_{bb}}{T_{bb}} \approx -2 \int \frac{k T_e}{m_e c^2} n_e \sigma_T dl$$

One also know that a perturbation in the temperature of the CMB is going to result in a perturbation BB intensity. In R-J limit

$$I_\nu = \frac{2\nu^3}{c^2} T_b$$

$$\delta I_\nu = \frac{2\nu^3}{c^2} \delta T_b$$

$$\Rightarrow \delta T_{bb} = \frac{c^2}{2\nu^3} \delta I_\nu$$

$$\Rightarrow \frac{\delta T_{bb}}{T_{bb}} \propto \frac{\delta I_\nu}{I_\nu}$$

[Can be verified observationally]

In R-J limit:

$$\frac{\delta T_{bb}}{T_{bb}} \approx -2 \frac{k_0}{m_e c^2} \sigma_T \int n_e T_e dl \quad \left( \propto - \frac{\delta I_\nu(R-J)}{I_\nu(R-J)} \right)$$

For clusters:

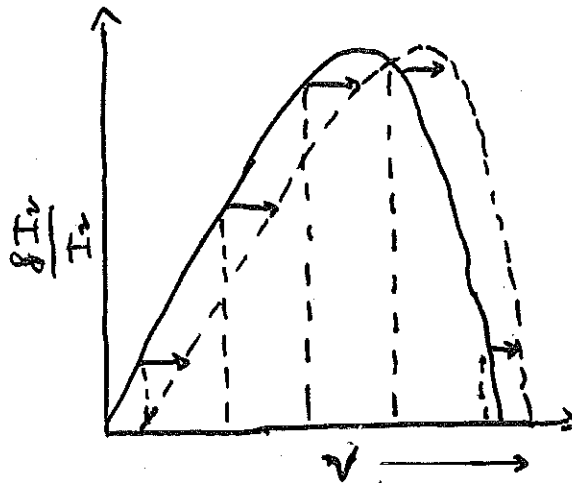
$$\left\{ \begin{array}{l} n_e \sim 10^{-2} \text{ cm}^{-3} \text{ and } L \sim \text{Mpc} \\ \frac{k T_e}{m_e c^2} \sim \frac{1}{50} \text{ with } k T_e \sim 10 \text{ keV} \end{array} \right.$$

In R-J limit:  $\frac{\delta T_{bb}}{T_{bb}} \approx -8 \times 10^{-4}$

$$T_{bb} \sim 3 \text{ K}$$

$$\delta T_{bb} \approx -10^{-3} \text{ K} \quad \text{[Significant depletion of CMB]}$$

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This is the result due to the IC scattering of the CMB photons off  $E_e \sim \text{few keV}$  electrons in clusters. IC drains photons from the CMB spectrum and duplicates the spectrum at higher frequencies.